

# Growth and spectroscopy of $\text{ZnWO}_4:\text{Ho}^{3+}$ crystal

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## Abstract

The single crystal of  $\text{ZnWO}_4:\text{Ho}^{3+}$  was grown by Czochralski technique. The XRD, absorption spectra, fluorescence spectrum are presented, and the Judd–Ofelt (J–O) intensity parameters  $\Omega_2$ ,  $\Omega_4$ ,  $\Omega_6$  are obtained to be  $4.61 \times 10^{-20} \text{ cm}^2$ ,  $0.16 \times 10^{-20} \text{ cm}^2$ ,  $0.48 \times 10^{-20} \text{ cm}^2$ , respectively. Calculated radiative transition rate, branching ratios, and radiative lifetime for different transition levels of  $\text{ZnWO}_4:\text{Ho}^{3+}$  crystals are presented. The most intense line correlative with the transition  $^5\text{S}_2 \rightarrow ^5\text{I}_8$  at 543 nm is potentially used for green solid state laser, and the fluorescence lifetime of this energy level is 14.1  $\mu\text{s}$ .

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## 1. Introduction

Holmium ion is known as the activator to lase at 2  $\mu\text{m}$  and 0.55  $\mu\text{m}$ . The main interest in 2  $\mu\text{m}$  corresponding to transition  $^5\text{I}_7 \rightarrow ^5\text{I}_8$  is for the use as an eye-safe source in atmosphere, medicine, wind shear, laser radar . . . [1–4]. The green emission corresponding to transition  $^5\text{S}_2 \rightarrow ^5\text{I}_8$ , centered at 0.55  $\mu\text{m}$ , has been observed, and can be utilized as a visible solid-state laser [5,6]. The ground state of  $\text{Ho}^{3+}$  is  $^5\text{I}_8$ , and the excited states are  $^5\text{I}_7$ ,  $^5\text{I}_6$ ,  $^5\text{I}_5$ ,  $^5\text{I}_4$ ,  $^5\text{F}_5$ ,  $^5\text{F}_4$ ,  $^5\text{G}_6$ ,  $^5\text{G}_5$ ,  $^3\text{H}_6$  . . .

Crystal  $\text{ZnWO}_4$  has the monoclinic structure with  $P2_1/c$  space group. Its lattice parameters  $a = 0.469263(5) \text{ nm}$ ,  $b = 0.572129(7) \text{ nm}$ ,  $c = 0.492805(5) \text{ nm}$ , and  $\beta = 90.6321(9)^\circ$  [7]. In this work, single crystal of  $\text{ZnWO}_4:\text{Ho}^{3+}$  was grown by Czochralski technique, and the spectroscopy was presented.

## 2. Experimental procedures

Using pure ZnO (99.9%),  $\text{WO}_3$  (99.9%),  $\text{Ho}_2\text{O}_3$  (99.97%), single crystal  $\text{ZnWO}_4:\text{Ho}^{3+}$  was grown by the Czochralski technique. Here, the doped concentration of  $\text{Ho}^{3+}$  is 2 at%. The initial compounds were mixed in a carnelian bowl, and sintered for almost 3–5 days at 1010  $^\circ\text{C}$ . Then the charge was deposited in a Pt crucible bowl of  $\varnothing 55 \text{ mm} \times 30 \text{ mm}$ , and placed in the DJL-400 furnace. The crystal was grown with the Pt-wire rotating at a rate of 8–12 rpm, and at a pulling rate of 1.2 mm/h. When the procedure was over, the crystal was drawn out, and cooled down to room temperature at a rate of 10–25  $^\circ\text{C}/\text{h}$ . Finally, the single-crystals  $\text{ZnWO}_4:\text{Ho}^{3+}$  with size of  $\varnothing 20 \text{ mm} \times 25 \text{ mm}$  were obtained.

The X-ray powder diffraction investigations are carried out with CAD4 diffractometer, equipped with Cu  $K\alpha$  radiation ( $\lambda = 1.54056 \text{ \AA}$ ). The data are collected using Ni-filtered Cu-target tube at room temperature in the  $2\theta$  range from  $5^\circ$  to  $85^\circ$ . Fig. 1 shows the diffraction patterns, which is in good accordance with the standard JCPDS card (No. 73 0554) of  $\text{ZnWO}_4$  [8].

The crystal was oriented by YX-2 X-ray Crystal Orientation Unit produced by Dandong Radioactive Instrument Co. Ltd., and was cut into several samples. Samples used for spectroscopic measurements are optically polished to flat.

Fig. 2 shows the room temperature absorption spectra of crystal  $\text{ZnWO}_4:\text{Ho}^{3+}$  measured by Lambda900 spectrometer (Perkin-Elmer UV-VIS Spectrometer) along three crystallographic axes. The samples are flat with the heights of 1.68 mm, 1.62 mm, and 1.64 mm along  $a$ ,  $b$ ,  $c$  directions, respectively. The fluorescence spectrum recorded by FLS920 at room temperature excited by 445 nm, using a CW xenon lamp as the exciting source, is shown in Fig. 3. The

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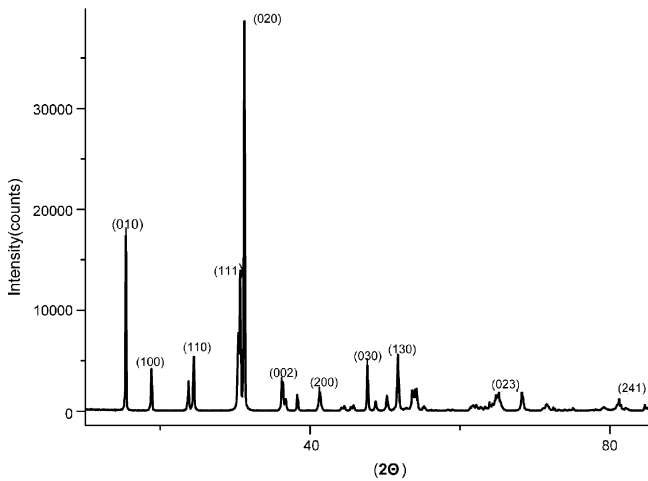


Fig. 1. X-ray powder diffraction pattern of  $\text{ZnWO}_4:\text{Ho}^{3+}$  crystal.

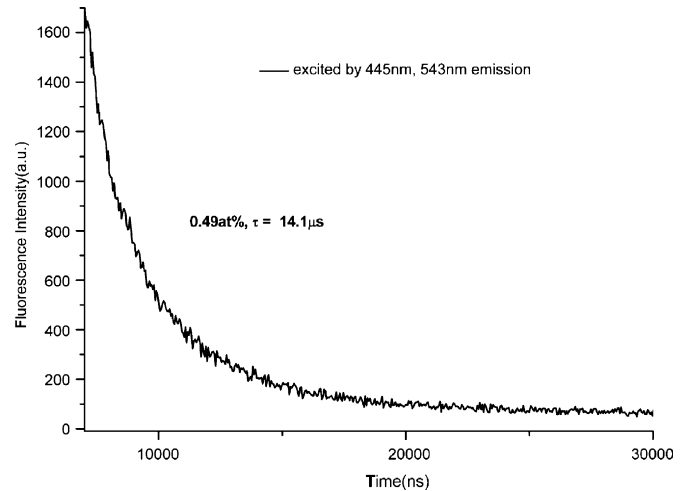


Fig. 4. Room temperature fluorescence decay curve excited by 445 nm, the y-axis is the fluorescence intensity, and the x-axis shows the time (ns).

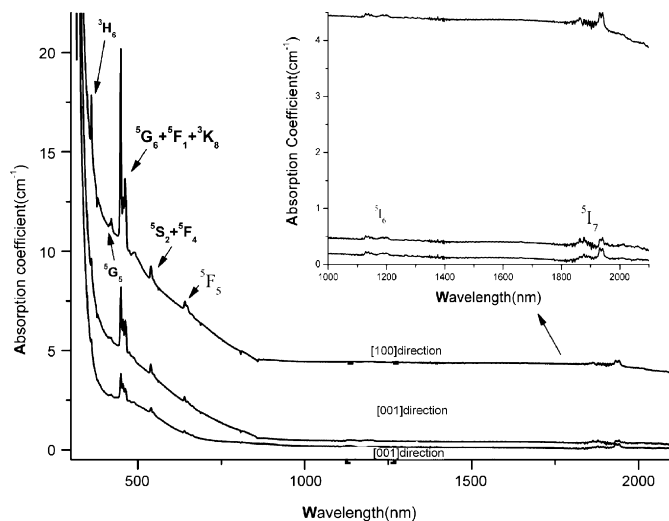


Fig. 2. Absorption spectra along  $a$ ,  $b$ ,  $c$  directions of crystal  $\text{ZnWO}_4:\text{Ho}^{3+}$ . The y-axis expresses absorption coefficient; the x-axis is the wavelength.

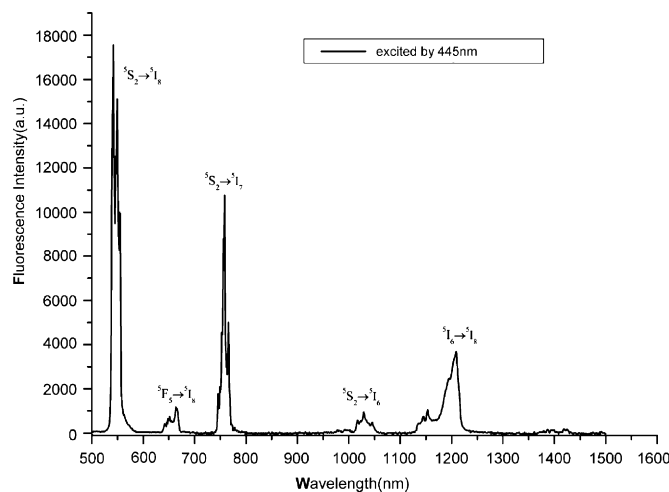


Fig. 3. Room temperature fluorescence spectrum of  $\text{ZnWO}_4:\text{Ho}^{3+}$  excited with 445 nm.

height of the used sample is 1.62 mm. The fluorescence decay lifetime curve correlative with  $^5\text{S}_2$  at 543 nm is shown in Fig. 4, and  $\tau_f$  is measured to be 14.1  $\mu\text{s}$ .

### 3. Results and discussion

The concentration of  $\text{Ho}^{3+}$  in  $\text{ZnWO}_4$  was measured to be 0.26 wt% (0.49 at%) by inductively coupled plasma atomic emission spectrometry method (ICP-AES instrument). So, the calculation of segregation coefficients of the dopant in crystals  $K$  was performed by the following formula:

$$K = \frac{(\text{molesHo}/(\text{molesHo} + \text{molesZnWO}_4))_{\text{crystal}}}{(\text{molesHo}/(\text{molesHo} + \text{molesZnWO}_4))_{\text{melt}}} \quad (1)$$

Here, molesHo, and moles $\text{ZnWO}_4$  are the mole fraction of dopant and trivalent host ion, respectively, in the crystal (numerator), and in the melt (denominator). The segregation coefficient of  $\text{Ho}^{3+}$  in  $\text{ZnWO}_4$  was obtained to be 0.24.

The monoclinic structure of  $\text{ZnWO}_4:\text{Ho}^{3+}$  crystal exhibits typical anisotropic characteristic. Thus, it is important to investigate the spectroscopy along three crystallographic axes. Fig. 2 shows the absorption spectra along [1 0 0] [0 1 0] [0 0 1] at room temperature. Along [1 0 0] direction, the base line moves away. The phenomenon could be explained by the natural [1 0 0] slip face in the structure of  $\text{ZnWO}_4$ . The absorption peaks are almost in the region of 300–800 nm. Five apparent transition peaks centered at 361 nm, 421 nm, 445 nm, 538 nm, and 641 nm are correlative with the transitions of  $^5\text{I}_8 \rightarrow ^3\text{H}_6$ ,  $^5\text{I}_8 \rightarrow ^5\text{G}_5$ ,  $^5\text{I}_8 \rightarrow ^5\text{G}_6 + ^5\text{F}_1 + ^3\text{K}_8$ ,  $^5\text{I}_8 \rightarrow ^5\text{F}_4 + ^5\text{S}_2$ , and  $^5\text{I}_8 \rightarrow ^5\text{F}_5$ , respectively. The transition of  $^5\text{I}_8 \rightarrow ^5\text{G}_6 + ^5\text{F}_1 + ^3\text{K}_8$  centered at 445 nm could be bumped to emit 0.55  $\mu\text{m}$  green laser. Zooming in the region of the infrared band, two peaks centered at 1128 nm and 1929 nm are found, which correlative with the transition of  $^5\text{I}_8 \rightarrow ^5\text{I}_6$  and  $^5\text{I}_8 \rightarrow ^5\text{I}_7$ . The transition of  $^5\text{I}_8 \rightarrow ^5\text{I}_7$  could be bumped to emit 2.0  $\mu\text{m}$  eye-safe laser. The absorption cross-section is calculated by the following equations:

$$\sigma_a = \frac{a}{N_c} \quad (2)$$

Table 1  
Integrated absorbance  $\Gamma$  along three crystallographic axes of  $\text{ZnWO}_4:\text{Ho}^{3+}$

Excited state (ground state $^5\text{I}_8$ )	Wavelength (nm)	$\Gamma$ (nm/cm)			$\overline{\Gamma(\lambda)}$ (nm/cm)
		<i>a</i>	<i>b</i>	<i>c</i>	
$^3\text{H}_6$	361	5.52	0.22	0.63	2.12
$^5\text{G}_5$	421	1.64	0.53	0.22	0.79
$^3\text{K}_8 + ^5\text{F}_1 + ^5\text{G}_6$	445	64.04	12.02	26.67	34.24
$^5\text{S}_2 + ^5\text{F}_4$	538	2.56	1.13	2.19	1.96
$^5\text{F}_5$	641	2.46	1.03	1.39	1.63
$^5\text{I}_6$	1128	3.27	3.20	3.38	3.28
$^5\text{I}_7$	1929	30.99	12.13	16.49	19.87

$$a = \frac{A}{L \log e} \quad (3)$$

Here,  $\sigma_a$  is the absorption cross-section;  $a$  is the absorption coefficient,  $A$  is the absorbance,  $L$  is the thickness of the polished crystal, and  $N_c$  is the doped ion concentration (here is 0.49 at%) in atoms. Based on the absorption spectra, the absorption cross-section versus wavelength was obtained. At 445 nm, the absorption cross-sections  $\sigma_a$  along  $a$ ,  $b$ ,  $c$  axes are  $15.5 \times 10^{-20} \text{ cm}^2$ ,  $3.7 \times 10^{-20} \text{ cm}^2$ , and  $7.3 \times 10^{-20} \text{ cm}^2$ , respectively. The average is  $8.8 \times 10^{-20} \text{ cm}^2$ . The absorption cross-sections centered at 1929 nm, correlative with transition  $^5\text{I}_8 \rightarrow ^5\text{I}_7$  and are  $6.0 \times 10^{-20} \text{ cm}^2$ ,  $0.34 \times 10^{-20} \text{ cm}^2$ , and  $0.55 \times 10^{-20} \text{ cm}^2$ , respectively.

The experimental data obtained from the absorption spectra are used to calculate the oscillator strengths. The experimental oscillator strength  $f_{\text{exp}}$  can be calculated by using the following formula:

$$f_{\text{exp}} = \frac{mc^2\Gamma}{\pi e^2 N_0 L \lambda^2} \quad (4)$$

$$\Gamma = \frac{\int D(\lambda) d\lambda}{L \log e} = \frac{2.303 \int D(\lambda) d\lambda}{L} \quad (5)$$

where  $N_0$  is the concentration of  $\text{Ho}^{3+}$ ;  $L$  is the thickness;  $n$  is the refractive index to be 2.25;  $\Gamma$  is the integrated absorbance for each absorption band;  $D(\lambda)$  is the absorbance. The integrated absorbance  $\Gamma$  associated with the seven transitions was calculated as shown in Table 1.

Based on the Judd–Ofelt theory, the measured line strengths are then used to obtain the J–O intensity parameters  $\Omega_2$ ,  $\Omega_4$ , and  $\Omega_6$  by fitting the set of equations from the corresponding transitions between J and J' manifolds in the following equation:

$$S_{\text{meas}}(J \rightarrow J') = \frac{(2J+1)}{N_0} \frac{3hc}{8\pi^3 \lambda e^2} \frac{9n}{(n^2+2)^2} \bar{\Gamma} \quad (6)$$

Table 2  
The J–O intensity parameters of  $\text{Ho}^{3+}$ -doped crystals

Crystal	$\Omega_2$ ( $10^{-20} \text{ cm}^2$ )	$\Omega_4$ ( $10^{-20} \text{ cm}^2$ )	$\Omega_6$ ( $10^{-20} \text{ cm}^2$ )	Reference
LaAlO <sub>3</sub>	5.362	1.862	1.065	[10]
YAlO <sub>3</sub>	1.82	2.38	1.53	[11]
YAG	0.04	2.67	1.89	[12]
LuAG	0.172	2.08	1.92	[13]
ZnWO <sub>4</sub>	4.61	0.16	0.48	This work

$$S_{\text{calc}}(J \rightarrow J') = \sum_{t=2,4,6} \Omega_t \left| \left\langle (S, L)J \left\| U^{(t)} \right\| (S', L')J' \right\rangle \right|^2 \quad (7)$$

$$S_{\text{calc}} = U\Omega \quad (8)$$

where  $U^{(t)}$  ( $t=2, 4, 6$ ) are the matrix elements of unit tensor that have been calculated by Carnall et al. [9]. The root mean square (rms) deviation between experimental, and calculated line strengths is determined by

$$\text{rms}\Delta S = \sqrt{\sum_{t=1}^N (S_{\text{meas}} - S_{\text{calc}})^2 / (N-3)} \quad (9)$$

where

$$\text{rms}S = \sqrt{\sum_{i=1}^N S_{\text{meas}}^2 / N} \quad (10)$$

So, the rms error is  $0.25 \times 10^{-20} \text{ cm}^2$ .

After a least-square fitting of  $S_{\text{meas}}$  to  $S_{\text{calc}}$ , the three J–O intensity parameters were obtained:  $\Omega_2 = 4.61 \times 10^{-20} \text{ cm}^2$ ,  $\Omega_4 = 0.16 \times 10^{-20} \text{ cm}^2$ , and  $\Omega_6 = 0.48 \times 10^{-20} \text{ cm}^2$ . The three J–O intensity parameters compared with other  $\text{Ho}^{3+}$ -doped crystals are shown in Table 2, [10–13]. The values of J–O parameters depend on the crystal structure. From the comparison, we can find that the values of  $\Omega_4$  and  $\Omega_6$  are smaller, which indicates the component of covalent bond is also smaller.

From the 4f–4f intensity model, the calculated oscillator strength of a transition between two multiplets is defined as [14]:

$$f_{\text{cal}} = \frac{8\pi^2 mc}{3h\lambda(2J+1)} \frac{(n^2+2)^2}{9n} \times \sum_{t=2,4,6} \Omega_t \left\langle 4f^n(\alpha' S' L')J' \left\| U^{(t)} \right\| 4f^n(\alpha S L)J \right\rangle^2 \quad (11)$$

Table 3  
Electrical dipole line strengths and oscillator strengths of ZnWO<sub>4</sub>:Ho<sup>3+</sup> crystal

Excited states	$\lambda$ (nm)	Line strength ( $10^{-20}$ cm <sup>2</sup> )		Oscillator strength ( $10^{-6}$ )	
		$S_{\text{exp}}$	$S_{\text{cal}}$	$f_{\text{exp}}$	$f_{\text{cal}}$
<sup>5</sup> I <sub>7</sub>	1929	1.005	0.862	0.819	0.702
<sup>5</sup> I <sub>6</sub>	1128	0.284	0.375	0.395	0.522
<sup>5</sup> F <sub>5</sub>	641	0.248	0.338	0.608	0.829
<sup>5</sup> S <sub>2</sub> + <sup>5</sup> F <sub>4</sub>	538	0.355	0.483	1.038	1.411
<sup>3</sup> K <sub>8</sub> + <sup>5</sup> F <sub>1</sub> + <sup>5</sup> G <sub>6</sub>	445	7.437	7.377	26.035	25.826
<sup>5</sup> G <sub>5</sub>	421	0.183	0.085	0.683	0.317
<sup>3</sup> H <sub>6</sub>	361	0.573	1.012	2.494	4.408

Table 4  
Calculated radiative transition rate, branching ratios, and radiative lifetime for different transition levels of ZnWO<sub>4</sub>:Ho<sup>3+</sup> crystal

Start levels	Terminal levels	Wavelength (nm)	$A$ (s <sup>-1</sup> )	$\beta$	$\tau_r$ ( $\mu$ s)
<sup>5</sup> S <sub>2</sub>	<sup>5</sup> F <sub>5</sub>	2979	0.23	$1.08 \times 10^{-4}$	463
	<sup>5</sup> I <sub>6</sub>	1028	127.44	0.06	
	<sup>5</sup> I <sub>7</sub>	757	812.47	0.37	
	<sup>5</sup> I <sub>8</sub>	543	1220	0.56	
<sup>5</sup> F <sub>5</sub>	<sup>5</sup> I <sub>6</sub>	1569	64.61	0.05	780
	<sup>5</sup> I <sub>7</sub>	1014	267.78	0.21	
	<sup>5</sup> I <sub>8</sub>	664	948.25	0.74	
<sup>5</sup> I <sub>7</sub>	<sup>5</sup> I <sub>8</sub>	1920	73.20	1	13660

$$\text{rms } f = \sqrt{\sum_{i=1}^N |f_{\text{cal}} - f_{\text{exp}}|^2 / N} \quad (12)$$

Then, the experimental oscillator strength  $f_{\text{exp}}$ , and calculated oscillator strength  $f_{\text{cal}}$  were obtained. The error  $\text{rms } f$  is obtained to be  $0.7 \times 10^{-6}$ . Table 3 lists the experimental, and calculated line strength  $S_{\text{meas}}$ ,  $S_{\text{calc}}$ , the oscillator  $f_{\text{exp}}$ , and  $f_{\text{calc}}$ .

The fluorescence spectrum of the crystal excited by 445 nm recorded at room temperature is shown in Fig. 3. In the spectra, five intense emission peaks at 543 nm, 664 nm, 757 nm, 1028 nm, and 1208 nm are corresponding to the transitions <sup>5</sup>S<sub>2</sub> → <sup>5</sup>I<sub>8</sub>, <sup>5</sup>F<sub>5</sub> → <sup>5</sup>I<sub>8</sub>, <sup>5</sup>S<sub>2</sub> → <sup>5</sup>I<sub>7</sub>, <sup>5</sup>S<sub>2</sub> → <sup>5</sup>I<sub>6</sub>, <sup>5</sup>I<sub>6</sub> → <sup>5</sup>I<sub>8</sub> of Ho<sup>3+</sup> ions, respectively. The most intense line at 543 nm with FWHM of 13 nm is potentially used for green laser.

Using the obtained emission line strengths, the radiative decay rates  $A(J \rightarrow J')$  can be determined by the following equations [15]:

$$A(J \rightarrow J') = \frac{64\pi^4 e^2}{3h(2J+1)\lambda^3} \frac{n(n^2+2)^2}{9} \times \sum_{t=2,4,6} \Omega_t \left| \langle (S, L)J \parallel U^{(t)} \parallel (S', L')J' \rangle \right|^2 \quad (13)$$

$$A_T(J) = \sum_{J'} A(J \rightarrow J') \quad (14)$$

Then the radiative lifetimes  $\tau_r = 1/A_T(J)$ , the mathematical formula for the fluorescent branching ratio is found by:

$$\beta(J') = \frac{A(J \rightarrow J')}{A_T(J)} \quad (15)$$

The calculated radiative transition rate, the branching ratios, and the radiative lifetime for different transition levels are presented in Table 4. The stimulated emission cross-section  $\sigma_p$  related to the radiative transition probability can be defined as:

$$\sigma_p = \frac{A(J \rightarrow J') \lambda_p^2}{4\pi^2 n^2 c \Delta \nu} \quad (16)$$

Here, the full frequency width at half maximum is  $\Delta \nu$ ;  $\lambda_p$  is the vacuum wavelength of the emission peak. Therefore, the emission cross-section  $\sigma_p$  centered at 543 nm is calculated to be  $0.14 \times 10^{-20}$  cm<sup>2</sup>.

Fig. 4 shows the fluorescence decay curve excited by 445 nm at room temperature in correspondence with the emission line <sup>5</sup>S<sub>2</sub> → <sup>5</sup>I<sub>8</sub> at 543 nm. The fluorescence lifetime  $\tau_f$  was measured to be 14.1  $\mu$ s, and the radiative lifetime  $\tau_r$  of the <sup>5</sup>S<sub>2</sub> level was calculated to be 463  $\mu$ s. The decay lifetime curve is almost single exponential, and apparently shorter than the radiative transitions calculated from the J–O theory. The phenomenon is probably from intense cross-relaxations down and up conversions  $\sigma_p \tau_f$  at 543 nm correlative with transition <sup>5</sup>S<sub>2</sub> → <sup>5</sup>I<sub>8</sub> is  $1.97 \times 10^{-26}$  cm<sup>2</sup> s.

#### 4. Conclusion

Single-crystal ZnWO<sub>4</sub>:Ho<sup>3+</sup> was grown by Czochralski technique. The absorption spectra along three crystallographic axes were measured at room temperature. Along three axes, anisotropic optical absorption characteristics are presented, and in the [1 0 0] direction, the baseline moves away due to the existence of slip faces in the structure of ZnWO<sub>4</sub>. Room temperature fluorescence spectrum was measured and discussed. The most

intense line at 543 nm corresponding to transition  $^5S_2 \rightarrow ^5I_8$  is found, whose emission cross-section is  $0.14 \times 10^{-20} \text{ cm}^2$ , and the fluorescence lifetime  $\tau_f$  of this energy level is measured to be 14.1  $\mu\text{s}$ .

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